

## Chemical Shift-correlated Two-dimensional Spin Echo N.M.R. Spectroscopy

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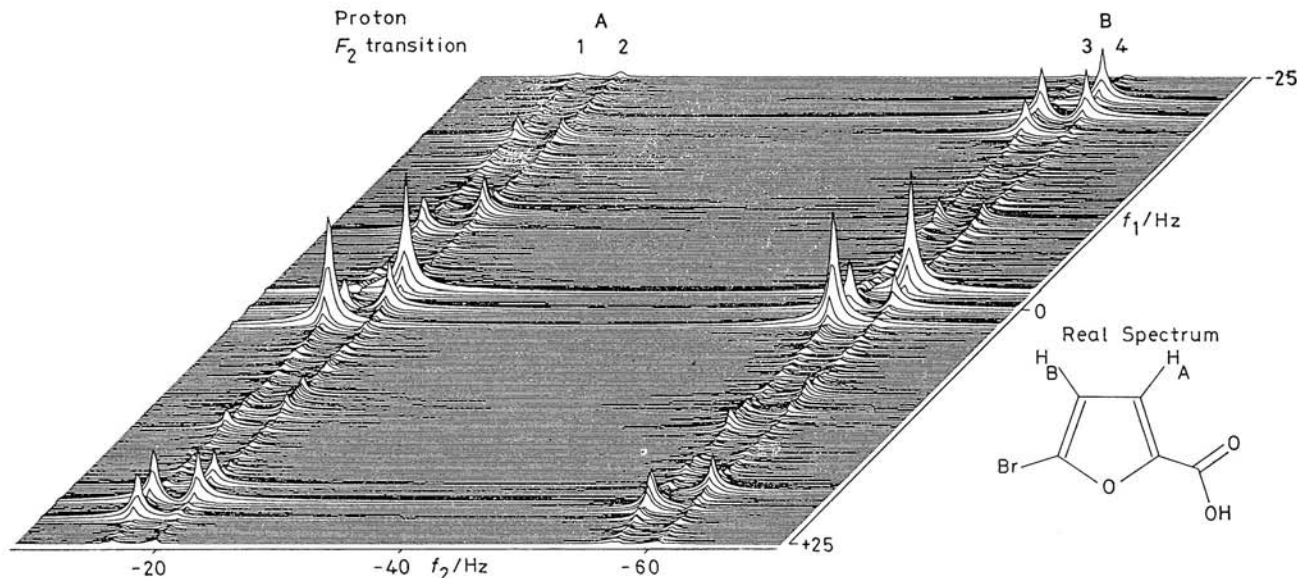
**Summary** An alternate two-dimensional n m r pulse sequence gives the chemical shifts of directly coupled nuclei

TWO-DIMENSIONAL (2D), homonuclear  $J$ -spectroscopy<sup>1-4</sup> has been applied to many complex problems in proton n m r spectroscopy<sup>5</sup> Free induction decays (F I D 's) with  $2m$  datum points are accumulated over an acquisition time  $t_2$  for a series of evolution times  $t_1$  in the pulse sequence  $90^\circ-t_1/2-180^\circ-t_1/2$  F I D Double Fourier transformation of the signal data matrix  $S(t_1, t_2)$  gives rise to a 2D spectrum in which signal intensity is plotted as a function of two frequency axes,  $f_1$  and  $f_2$  There is a  $J$ -spectrum, parallel to the  $f_1$  axis, for each of the  $m$  points along  $f_2$  The projection of the 2D spectrum on to the  $f_2$  axis gives the conventional 1D spectrum In the limit of weak coupling and perfect  $180^\circ$  pulses, the  $J$ -spectrum of any transition in  $f_2$  will contain only a single line

Much more information is gained from a  $90^\circ-t_1/2-90^\circ-t_1/2$ -F I D pulse sequence In general, the spectrum in  $f_1$  ( $J$ -spectrum) of a transition in  $f_2$  will contain many lines, owing to the mixing of that transition with every other transition of the same spin or of a directly coupled spin

Furthermore, the frequencies of  $f_1$  can arise from both the sum and difference of  $f_2$  frequencies The chemical shift of any one nucleus coupled to another may be found by inspection of the  $J$ -spectra of the second nucleus, even if all the transitions of the first nucleus are hidden in the 1D spectrum This technique is a development of a previous 2D experiment<sup>2,6</sup> in which the second delay time was omitted The second delay time ensures that at least some of the line broadening due to magnetic field inhomogeneities is eliminated Furthermore since the  $J$ -spectra start at  $f_1 = 0$  Hz a considerable saving in  $f_1$  digitization can be achieved over the original experiment,<sup>2</sup> as has recently been pointed out<sup>7</sup> The technique may be considered as an F1 alternative to the homonuclear INDOR experiment,<sup>8</sup> or selective decoupling

We report here the results of applying the  $90^\circ-t_1/2-90^\circ-t_1/2$ -F I D pulse sequence to the AB protons in 5-bromofuroic acid Figure 1 shows the 2D spectrum revealing many lines in  $f_1$  for each transition in  $f_2$  Figure 2(A) shows the  $J$ -spectrum of transition 2 [ $(\omega_A - J/2)$  in  $f_2$ ] Figure 2(B) shows a computer simulation of this  $J$ -spectrum produced on a 3-spin two-dimensional simulation program,<sup>‡</sup> based on the superspin formalism<sup>9</sup>



**FIGURE 1** The 2D proton n m r spectrum of 5-bromofuroic acid in  $\text{Me}_2\text{SO}$  at 90 MHz F I D 's were accumulated over 2048 datum points with  $t_1$  being incremented by 0.010 s 179 times, from an initial value of 0.001 s sweep width =  $\pm 125$  Hz delay between pulse sequences 2 s After the first Fourier transform was performed a fixed phase-correction was given to the spectra A magnitude calculation was applied after the second Fourier transformation No line-broadening was used

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‡ A listing of this program is available upon request to Dr A D Bain

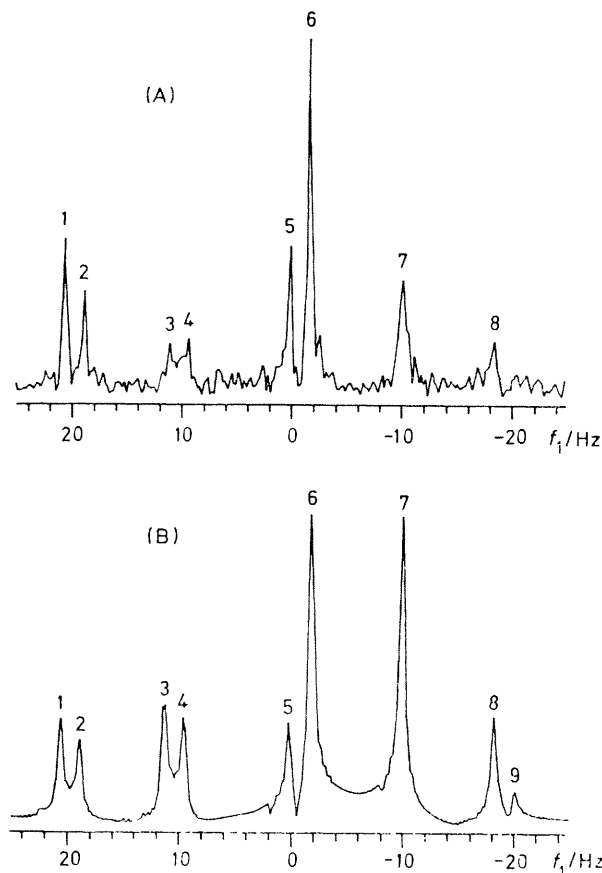


FIGURE 2. (A) The real  $J$ -spectrum of transition 2; (B) the simulated  $J$ -spectrum of transition number 2. The line intensities depend strongly on the flip angle of the refocussing pulse ( $120^\circ$  pulses were used in the simulation). Furthermore  $B_0$  inhomogeneity will have differential effects on line intensities.

The analysis of the  $J$ -spectra is quite simple. In this case, each transition in  $f_2$  can be mixed with every other one, since these transitions either belong to coupled spins or they belong to the same spin. Three distinct sets of peaks will arise in  $f_1$  for each transition in  $f_2$ . Firstly, subtractive

§ All frequencies in  $f_2$  are relative to the carrier frequency at 0 Hz.

¶ After submitting this communication a report on the application of the method to the analysis of the  $^1\text{H}$  n.m.r. spectrum of a small protein has appeared; see ref. 7.

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mixing of a transition with others in  $f_2$  will give a spectrum of four peaks starting at  $f_1 = 0$  Hz, with the transition mixing with itself [ $(\omega_i - \omega_i)/2 = 0$  Hz]. Secondly, additive mixing will give a spectrum of four peaks, starting at  $f_1 = \omega_i$  Hz ( $\omega_i$  is the appropriate  $f_2$  transition frequency)§ with a peak due to the mixing of a line with itself

TABLE. The predicted and observed frequencies of the peaks in the  $J$ -spectrum of the transition in  $f_2$  at  $\omega_A - J/2$  Hz (transition number 2)<sup>a</sup>

Peak No.	Origin	Predicted freq./Hz	Observed freq./Hz
5	$(\omega_A - J/2) - (\omega_A - J/2)$	0	+ 0.05
6	$(\omega_A - J/2) - (\omega_A + J/2)$	- 1.79	- 1.79
2	$(\omega_A - J/2) - (\omega_B + J/2)$	+18.70	+18.74
1	$(\omega_A - J/2) - (\omega_B - J/2)$	+20.49	+20.50
9	$(\omega_A - J/2) + (\omega_A - J/2)$	-20.03	<sup>c</sup>
8	$(\omega_A - J/2) + (\omega_A + J/2)$	-18.26	-18.48
3	$(\omega_A - J/2) + (\omega_B + J/2)$	+11.25 <sup>b</sup>	+10.94 <sup>b</sup>
4	$(\omega_A - J/2) + (\omega_B - J/2)$	+9.46 <sup>b</sup>	+9.24 <sup>b</sup>
7	$(\omega_A - J/2)$	-10.02	-10.16

<sup>a</sup>  $(\omega_A + J/2) = -16.48$ ,  $(\omega_A - J/2) = -20.03$ ,  $(\omega_B + J/2) = -57.43$ ,  $(\omega_B - J/2) = -61.02$  Hz. <sup>b</sup> Folded lines. <sup>c</sup> Not observed, but note particularly low intensity of this line in the simulation.

[ $(\omega_i + \omega_i)/2 = \omega_i$  Hz]. Thirdly, a phantom<sup>4</sup> peak will arise in  $f_1$  at  $\omega_i/2$  Hz. The Table lists the predicted and observed  $f_1$  frequencies of the peaks in transition number 2. It will be noted that all the frequencies in  $f_1$  are reduced by a factor of 2 after mixing,  $(\omega_i \pm \omega_i)/2$ , as in the  $90^\circ-t_1/2-180^\circ-t_1/2$ -echo experiment.<sup>1</sup> The analysis given here applies equally well to strongly or weakly coupled systems. For this experiment, the second pulse can only mix lines present in the normal 1D spectrum, or create a phantom.

This technique promises to be valuable for finding the chemical shifts of hidden, coupled spins in n.m.r. spectra. We have used this technique to solve the complicated assignment problems in the 90 MHz  $^1\text{H}$  n.m.r. spectra of some ribonucleic acids, and will report on this application shortly.¶

We thank the National Science and Engineering Research Council of Canada for financial assistance.

(Received, 17th September 1979; Com. 997.)